

Importance of Coherent Beam-Beam Effects to the Compensation of Beam- Beam Tune Spread in Hadron Collider

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Features of the Code

- Self-consistent beam-beam kicks in transverse space are calculated with PIC method.
 - Size of adaptive grid is matched with beam.
 - Large number of macro-particles is necessary for hadron beams in nonlinear regime of beam-beam interactions.
- Computing
 - dynamics of beam tune spread
 - dynamics of beam particle distributions



Formulas and method

- Kick at IPs:

$$\Delta \vec{r}' = \int d\vec{r} \rho(\vec{r}) \vec{G}(\vec{r} - \vec{r}) \quad \vec{r} = (x, y)$$

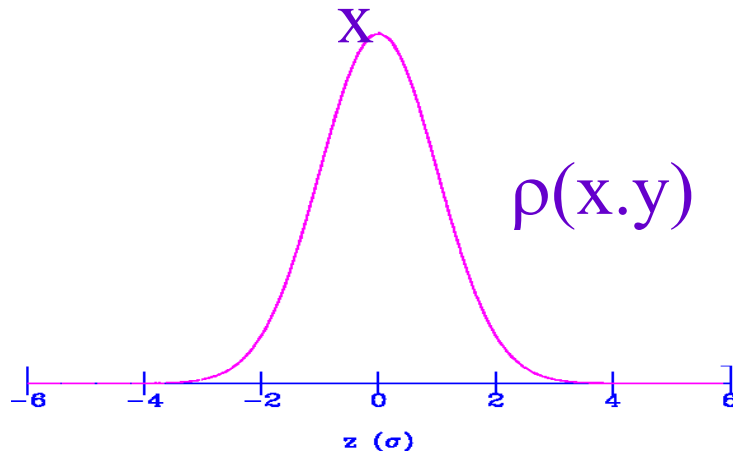
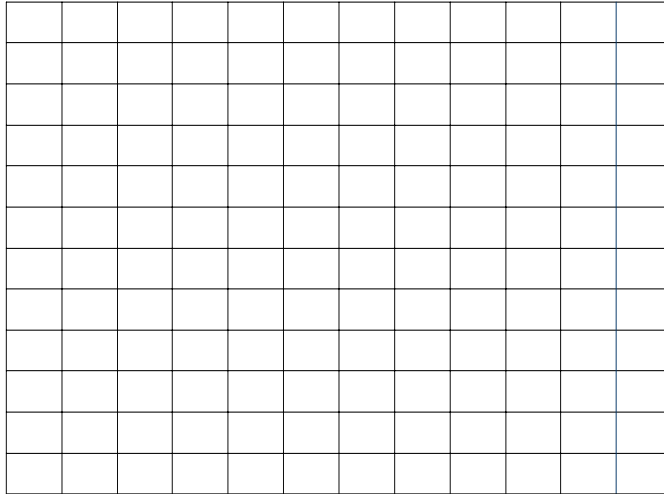
$$\vec{G}(\vec{r} - \vec{r}) = G_0 \frac{(\vec{r} - \vec{r})}{(\vec{r} - \vec{r})^2}$$

$$G_0 = 2Nr / \gamma = \begin{cases} 4\pi \xi_x \sigma_x^* (\sigma_x^* + \sigma_y^*) / \beta_x^* \\ 4\pi \xi_y \sigma_y^* (\sigma_x^* + \sigma_y^*) / \beta_y^* \end{cases}$$



PIC Calculation of Beam-Beam Kick

y



The x-y space covered by an uniform mesh

$\rho(x,y)$ obtained by 4-point-cloud in cell

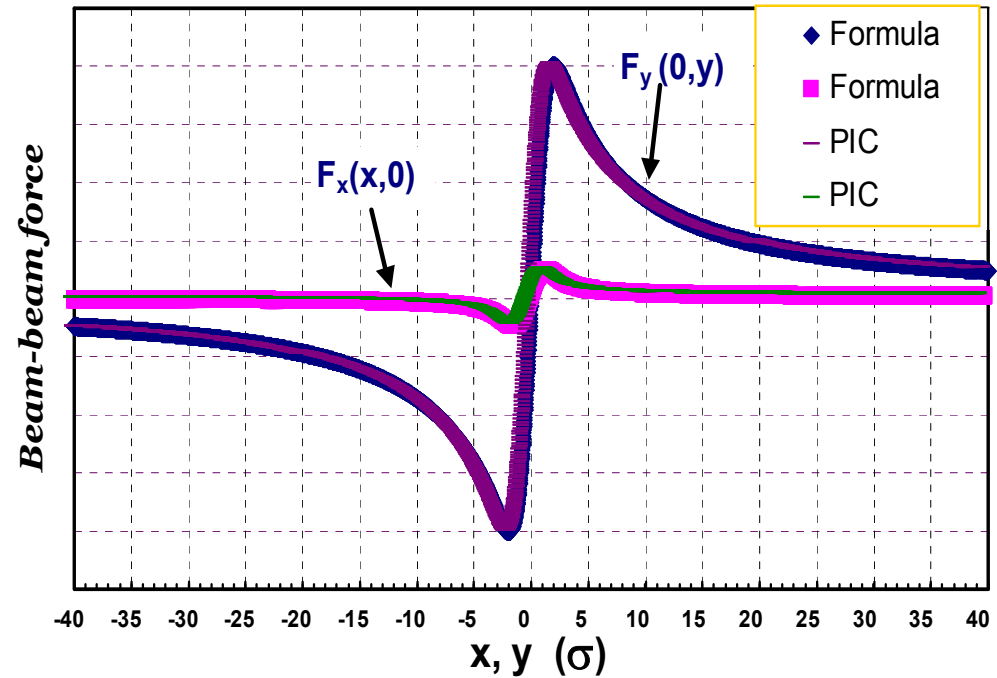
The forces calculated at the grid points

The fields interpolated to the position of every particle



Convergence of the Code

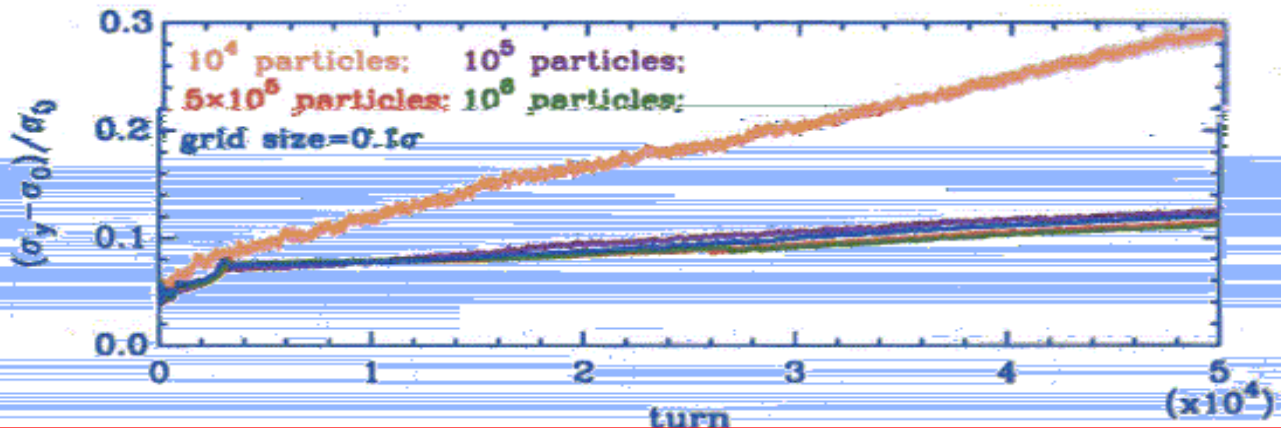
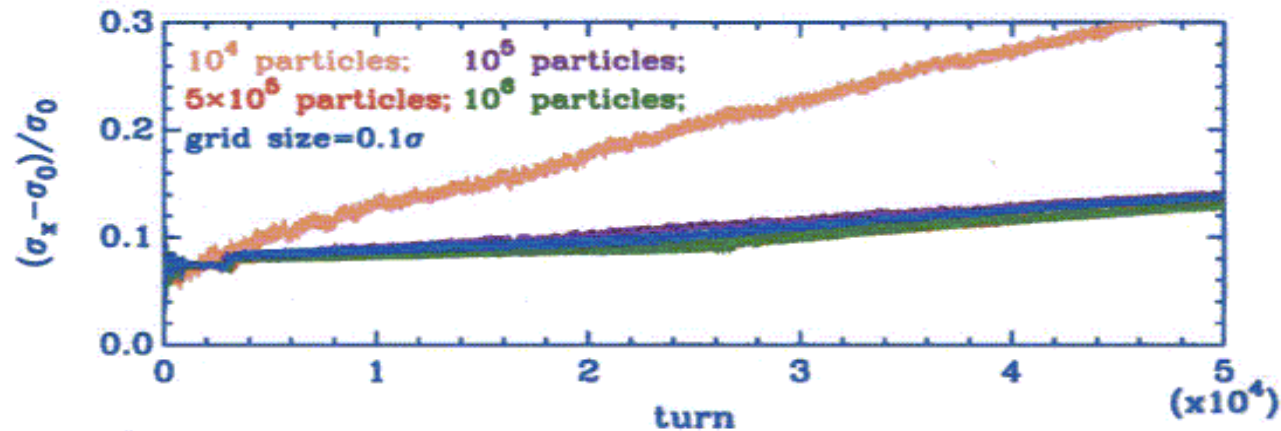
- number of macro-particles
- size of mesh
- grid constant





COMPARISON BETWEEN DIFFERENT NUMBERS OF MACRO-PARTICLES

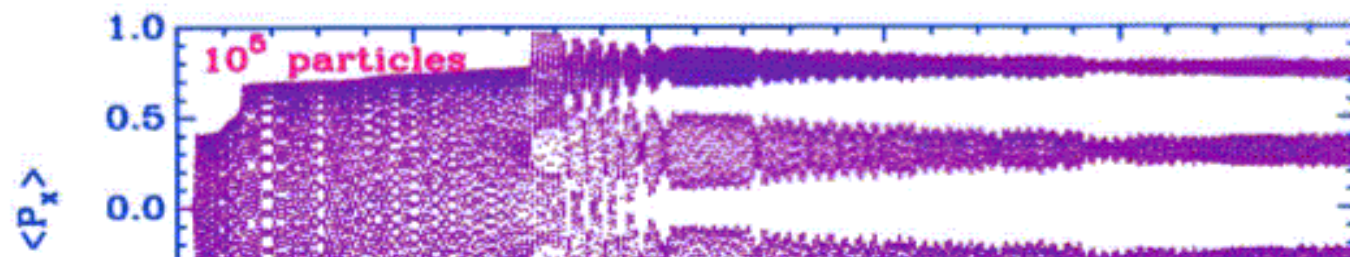
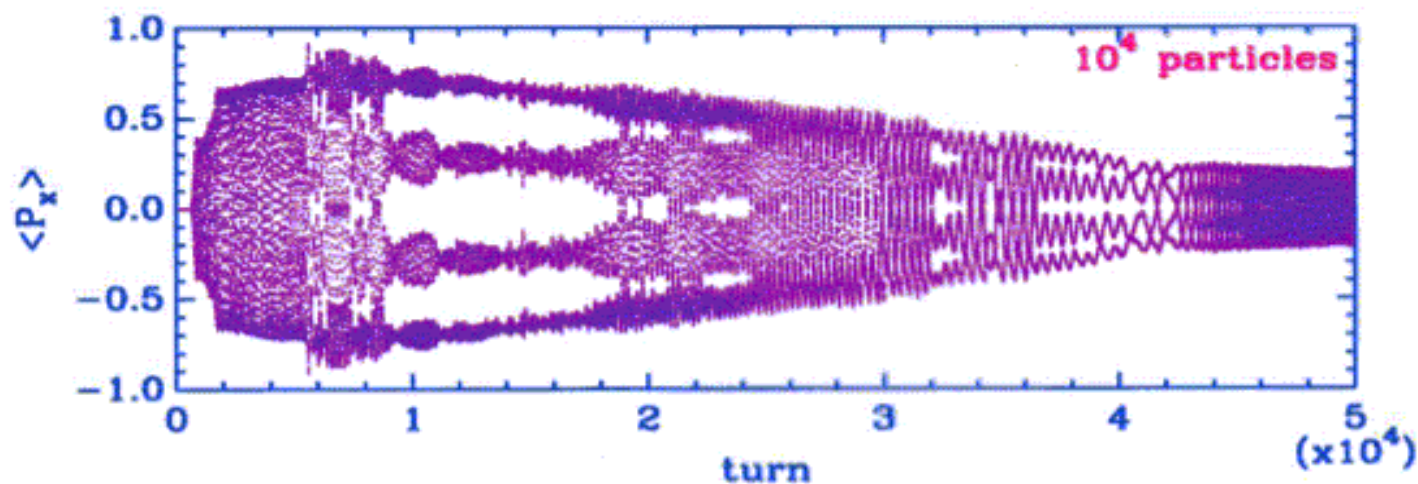
$$(\nu_x=0.31, \nu_y=0.32, \xi=0.04)$$



Fewer than 10^5 macro-particles are not enough to reveal true beam dynamics in nonlinear regime of beam-beam interactions.



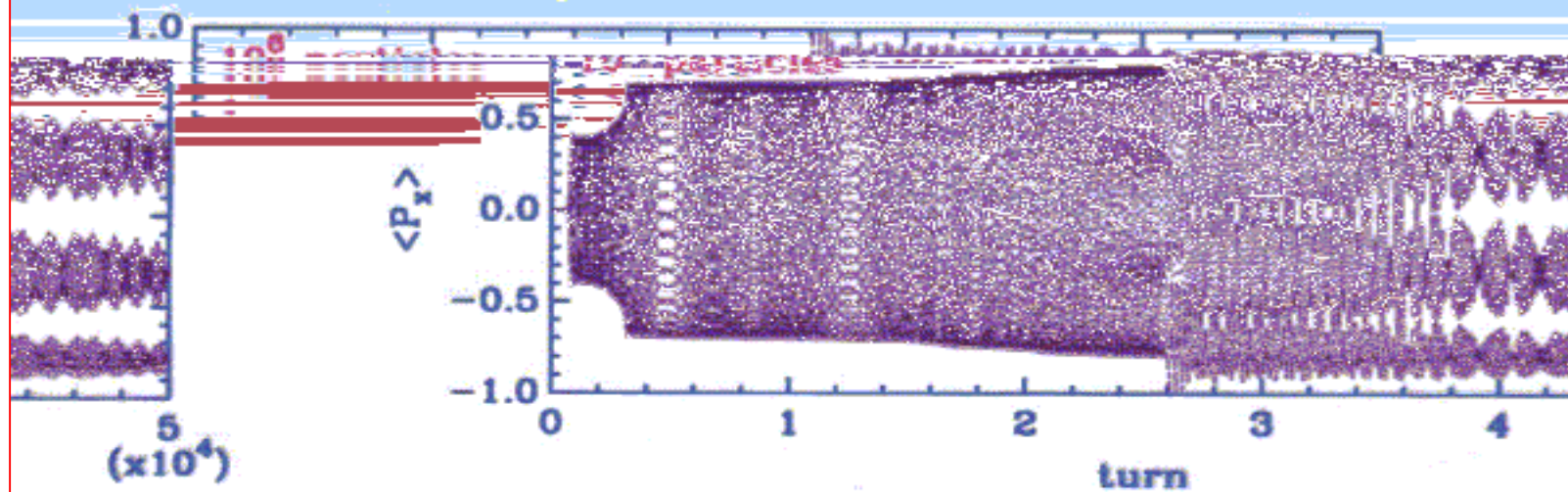
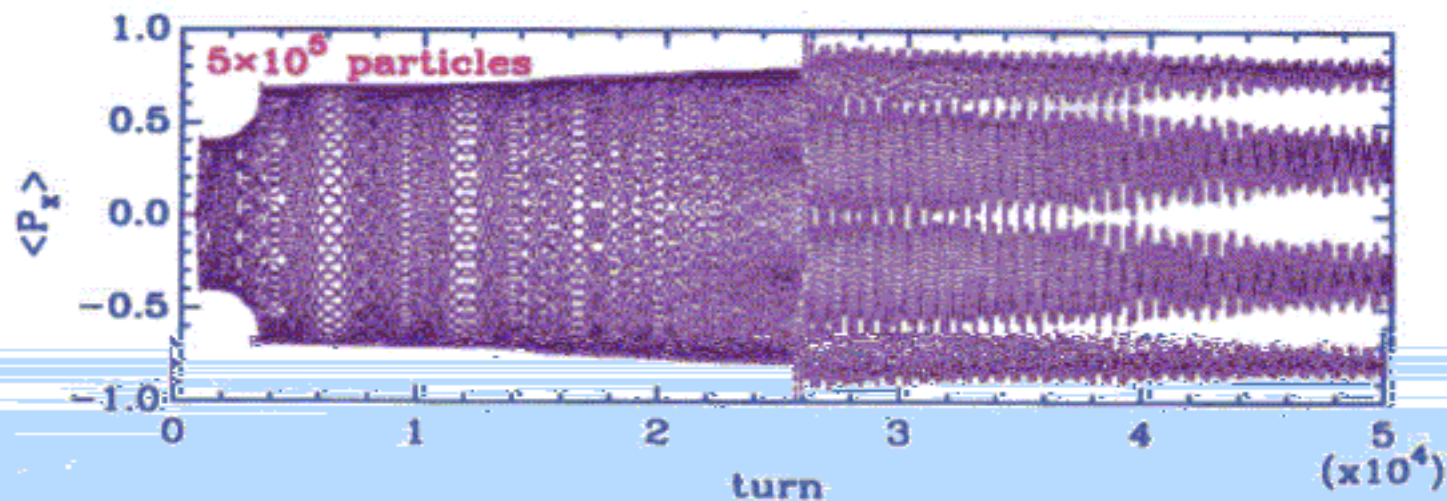
COMPARISON BETWEEN DIFFERENT NUMBERS OF MACRO-PARTICLES ($\nu_x=0.31$, $\nu_y=0.32$, $\xi=0.04$)





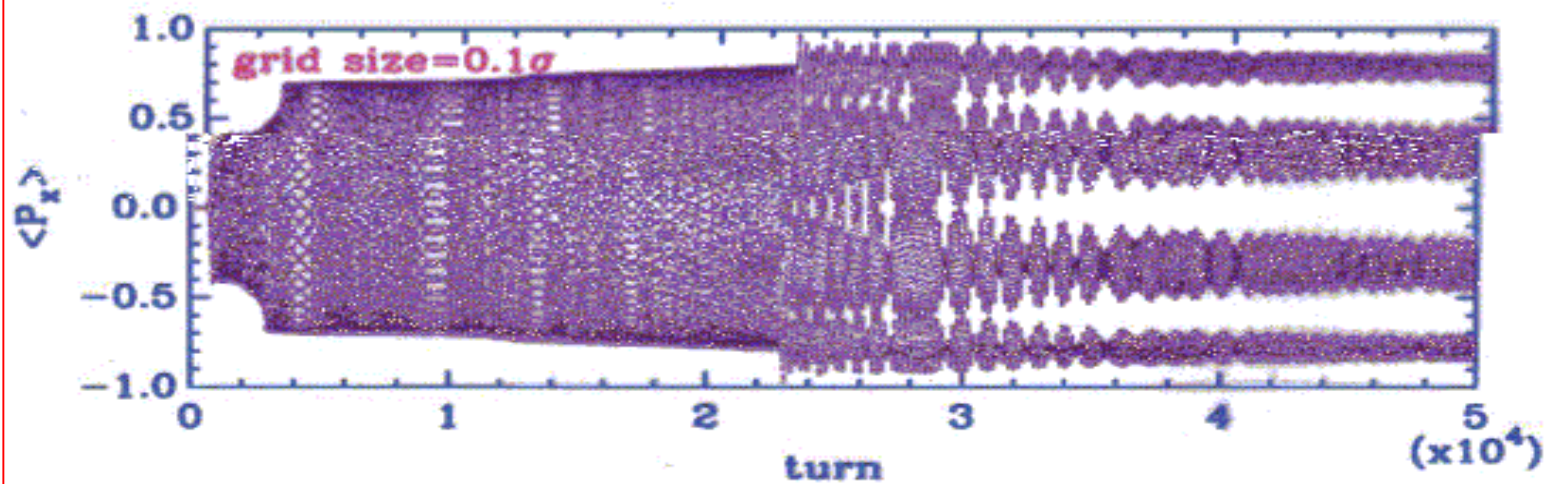
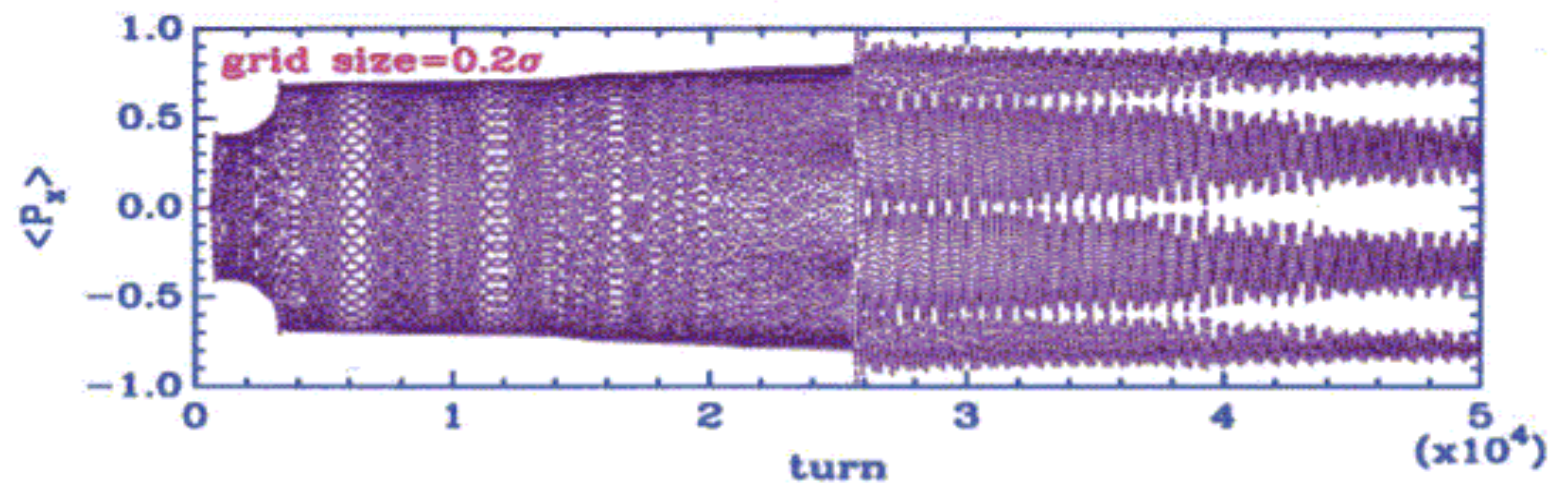
COMPARISON BETWEEN DIFFERENT NUMBERS OF MACRO-PARTICLES

($\nu_x=0.31$, $\nu_y=0.32$, $\xi=0.04$)





COMPARISON BETWEEN DIFFERENT GRID CONSTANTS ($\nu_x=0.31$, $\nu_y=0.32$, $\xi=0.04$)





INTRODUCTION

- ❏ A large beam-beam tune spread may lead to crossings of many high-order resonances. The compensation of the beam-beam tune spread has therefore been explored for a reduction of nonlinear beam-beam effects in hadron colliders.
- ❏ In this work, the importance of beam-beam tune spread to the chaotic coherent beam-beam instability was studied. It showed that in the nonlinear regime of beam-beam interactions, a beam-beam tune spread of a certain size is usually necessary to the stability of hadron beams.
- ❏ Recent beam experiments on HERA confirmed this conclusion.
- ❏ The study was conducted with a self-consistent beam-beam simulation with PIC method in model lattices of Tevatron and LHC.

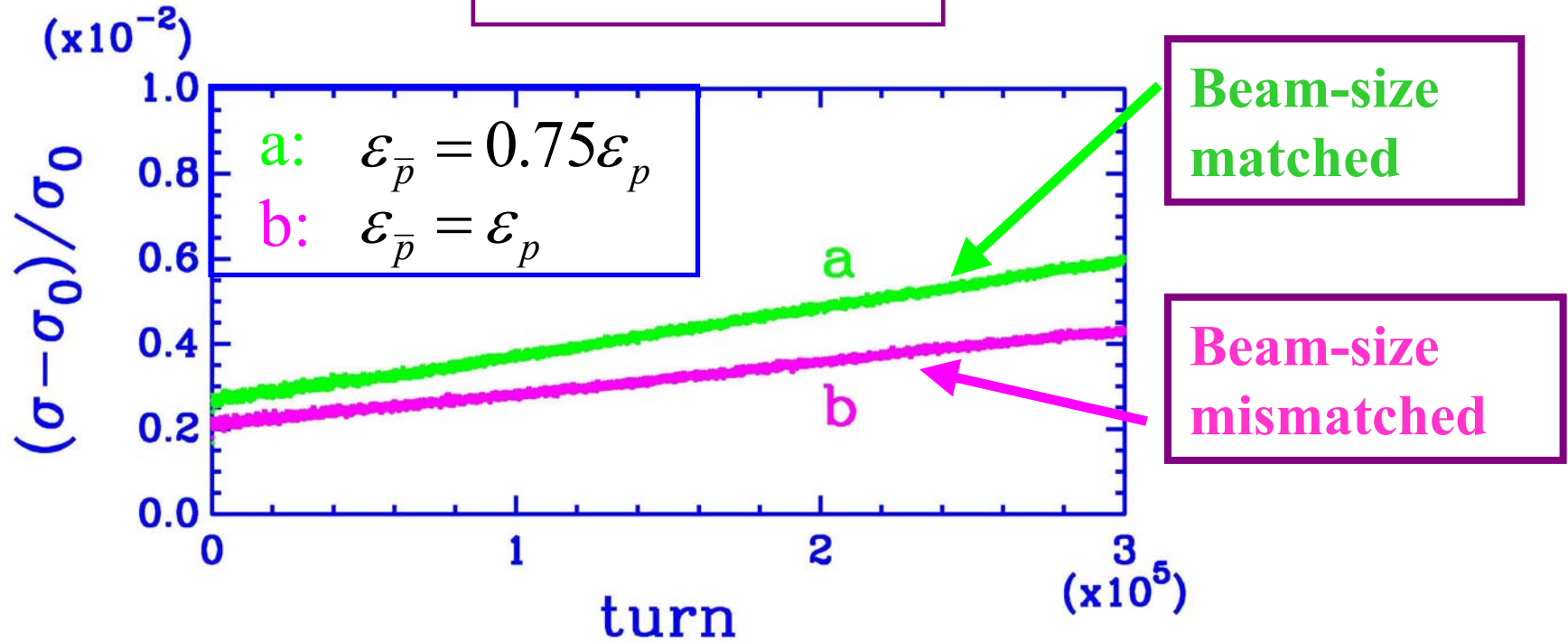


Evolution of r.m.s. Beam-Size of the Antiproton Beam

TEVATRON MODEL

$$(\nu_x, \nu_y) = (0.582, 0.574)$$

$$\xi_{\bar{p}p} = -0.01, \xi_{p\bar{p}} = -0.002$$



The beam-size growth rate in the mismatched collisions is larger than that in the matched case.



“Bad” and “Good” of Beam-Beam Tune Spread

$$H = v \square I + H_{\text{beam-beam}}(I, \phi, t)$$

$$= v \square I + \underbrace{\langle H_{\text{beam-beam}} \rangle}_{\substack{\text{Average part} \\ \text{integrable} \\ \text{“linear”}}} + \underbrace{\{H_{\text{beam-beam}}\}}_{\substack{\text{Oscillating part} \\ \text{nonintegrable} \\ \text{“nonlinear”}}}$$

In near-linear regime,
 $\langle H_{\text{beam-beam}} \rangle$ **dominates**
beam-beam interactions.

In nonlinear regime,
 $\{H_{\text{beam-beam}}\}$ **is important.**

Bad :

□ A large tune spread could result in crossings of bad resonance in case of a “bad” working point.

Good:

□ In nonlinear regime, a larger tune spread could result in a stronger Landau damping that could suppress chaotic coherent beam-beam instability.

□ Existence of a large tune spread reduces the possibility of trapping particles inside a resonance.

Comment:

For high-intensity hadron beams, beam-beam interaction is likely to be in the nonlinear regime and the beam-beam tune spread does more good than bad to the beam stability, except in the case of bad working point.



Compensation of Beam-Beam Tune Spread with Electron Beams

$$H = \vec{v} \cdot \vec{I} + H_{\bar{p}p}(I, \phi, t) + H_{\bar{p}e}(I, \phi, t)$$

$$= \vec{v} \cdot \vec{I} + (\langle H_{\bar{p}p} \rangle + \langle H_{\bar{p}e} \rangle) + (\{H_{\bar{p}p}\} + \{H_{\bar{p}e}\})$$

e-beam compensation:

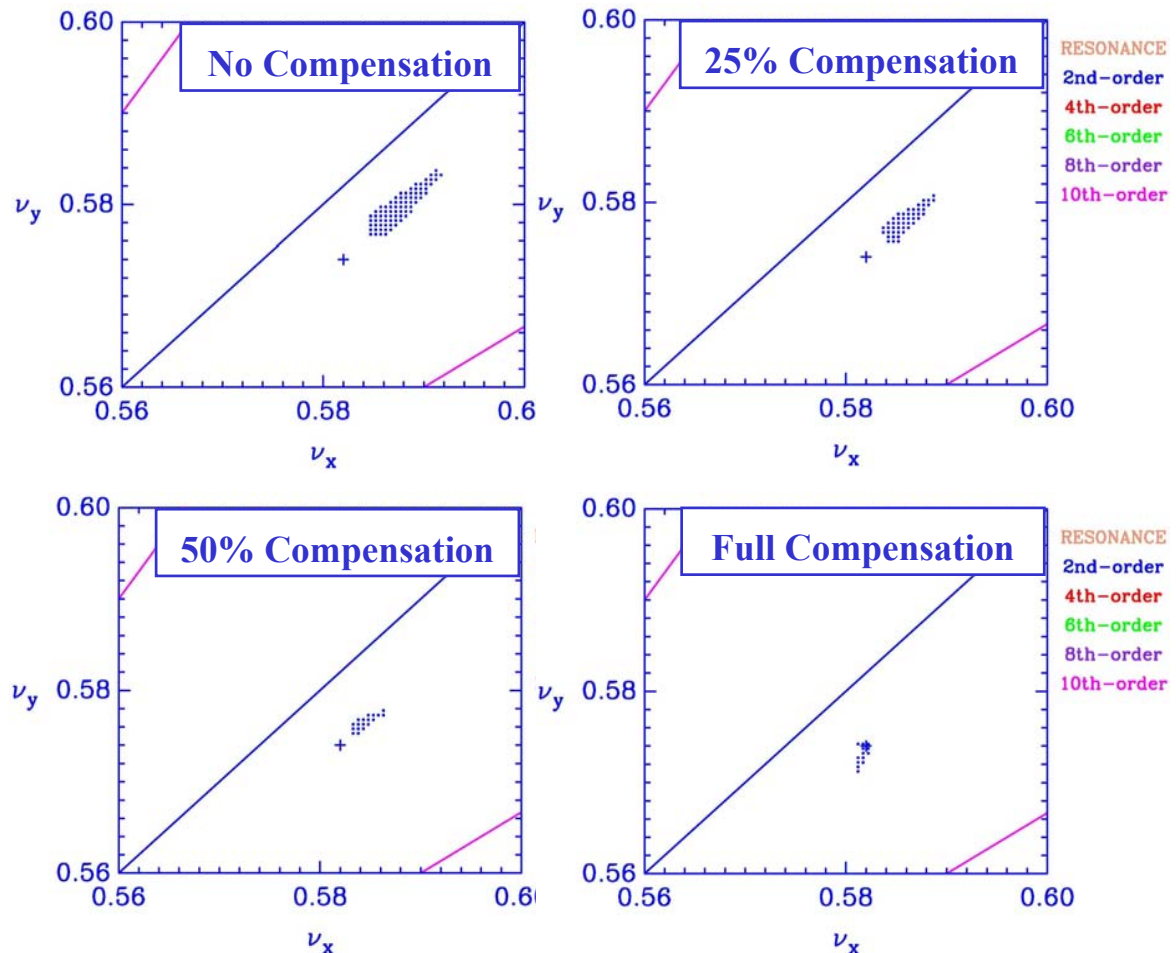
Reduce $(\langle H_{\bar{p}p} \rangle + \langle H_{\bar{p}e} \rangle)$

TEVATRON MODEL

$$(\nu_x, \nu_y) = (0.582, 0.574)$$

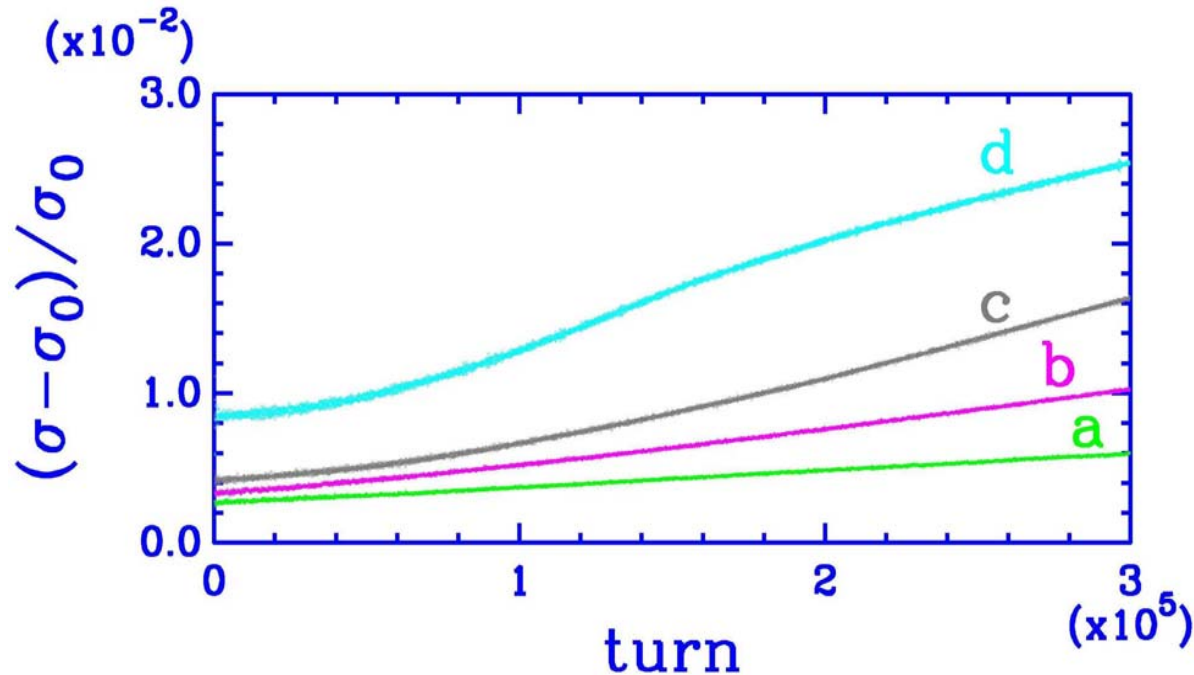
$$\xi_{\bar{p}p} = -0.01, \xi_{p\bar{p}} = -0.002$$

Tune Spread of the antiproton beam with or without the Compensation





Evolution of r.m.s. Beam-size of the Antiproton Beam



TEVATRON MODEL

$$(\nu_x, \nu_y) = (0.582, 0.574)$$

$$\xi_{\bar{p}p} = -0.01, \xi_{p\bar{p}} = -0.002$$

a : No compensation

b : 25% compensation

c : 50% compensation

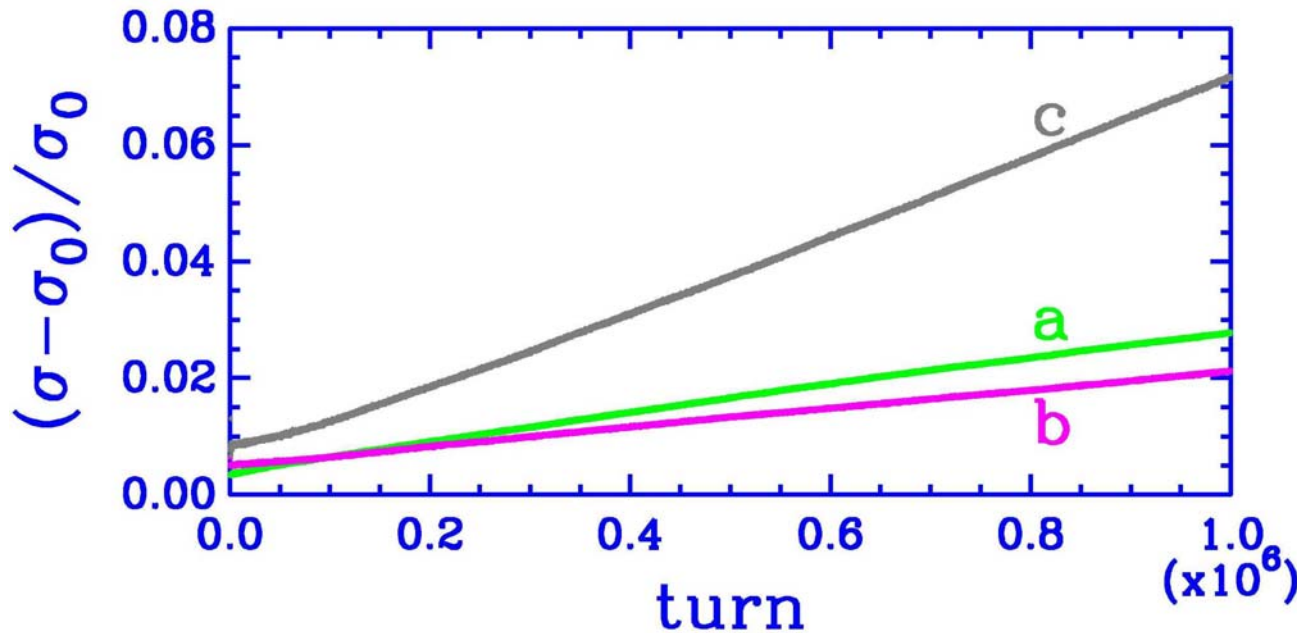
d : Full compensation



At the nominal working point, the compensation of beam-beam tune spread with electron beams could damage the beam stability .



Evolution of r.m.s. Beam-size of the Antiproton Beam When Two Beams Have Different Tunes



TEVATRON MODEL

$$\bar{p}: (\nu_x, \nu_y) = (0.582, 0.574)$$

$$p: (\nu_x, \nu_y) = (0.587, 0.579)$$

a : No compensation

b : 50% compensation

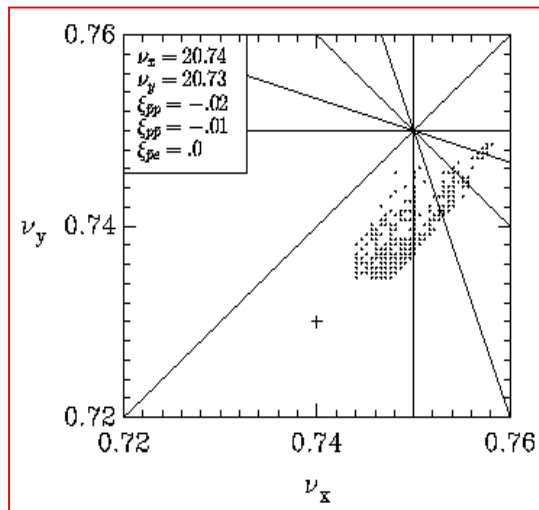
c : Full compensation

- When the tune split is big enough (0.005), the tune-spread compensation up to a certain degree could benefit the pbar beam.
- The fact that the tune of the p beam has an impact on the dynamics of the pbar beam confirms the existence of collective beam-beam effects in a strong-weak situation of beam-beam interactions.

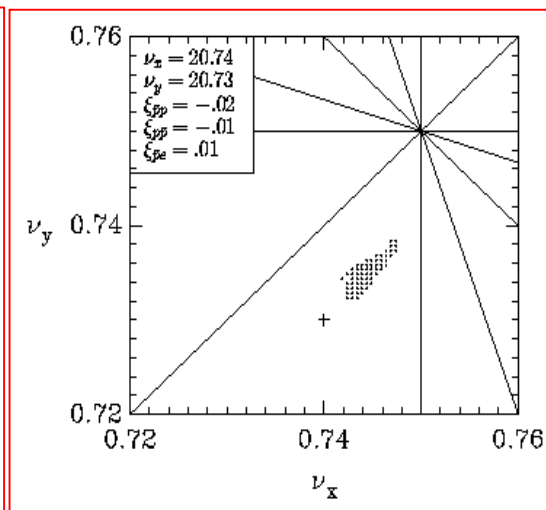


Tunes Close to the 4th-Order Resonance

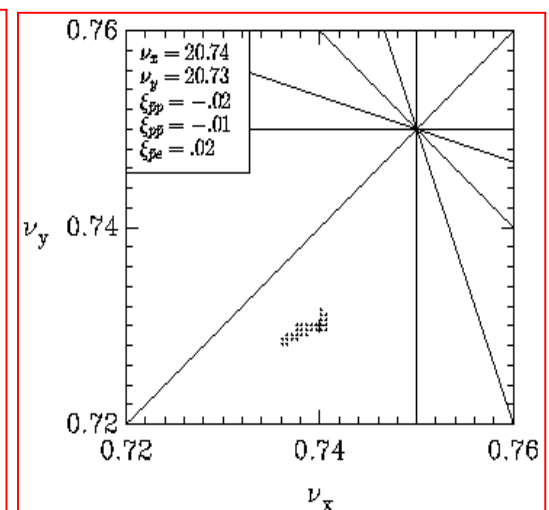
Beam-Beam Tune Spread of the Antiproton Beam



No compensation

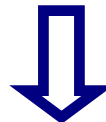


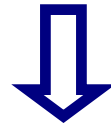
Half compensation



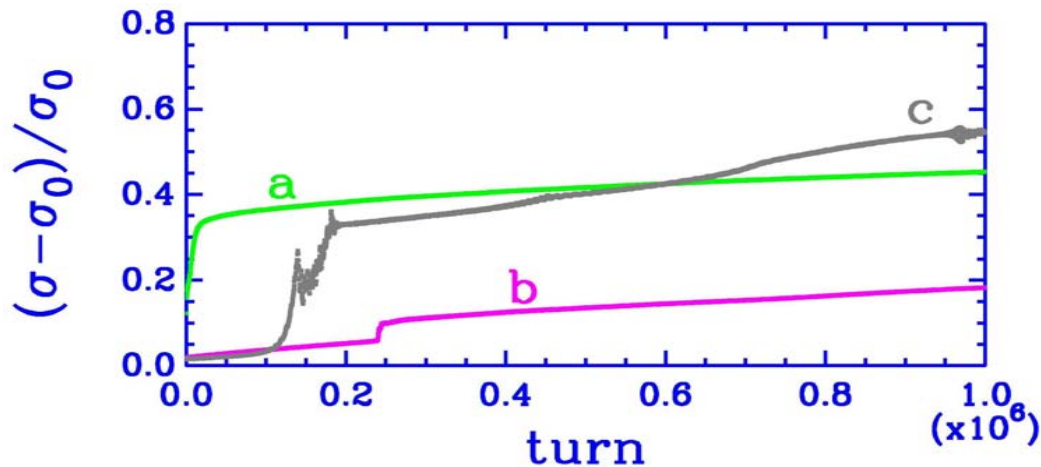
Full compensation

Crossing of a major resonance can be avoided with the tune-spread compensation .





Evolution of r.m.s. Beam-size of the Antiproton Beam



a: no compensation

b: 50% compensation

c: Full compensation

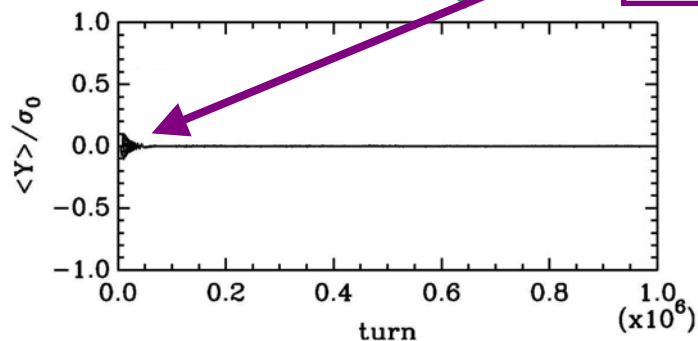
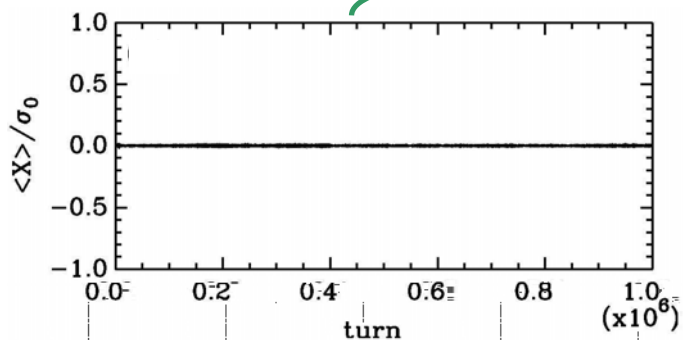
At a “bad” working point, a proper reduction of beam-beam tune spread could benefit beams.



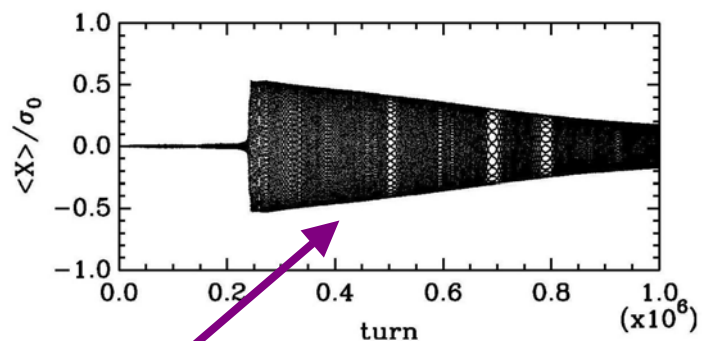


Beam-centroid Motion of the Antiproton Beam

No compensation

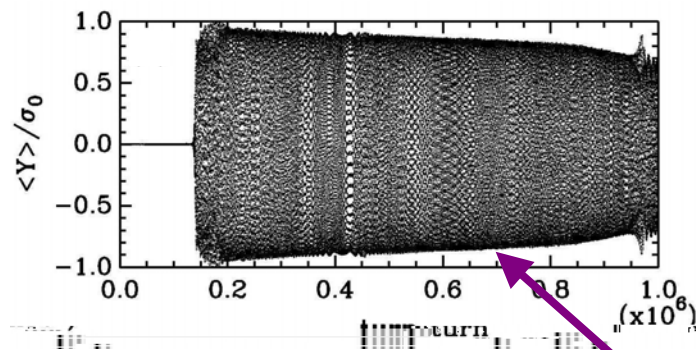


Very strong
Landau damping



50% compensation

Strong Landau
damping



Full compensation

Weak Landau
damping

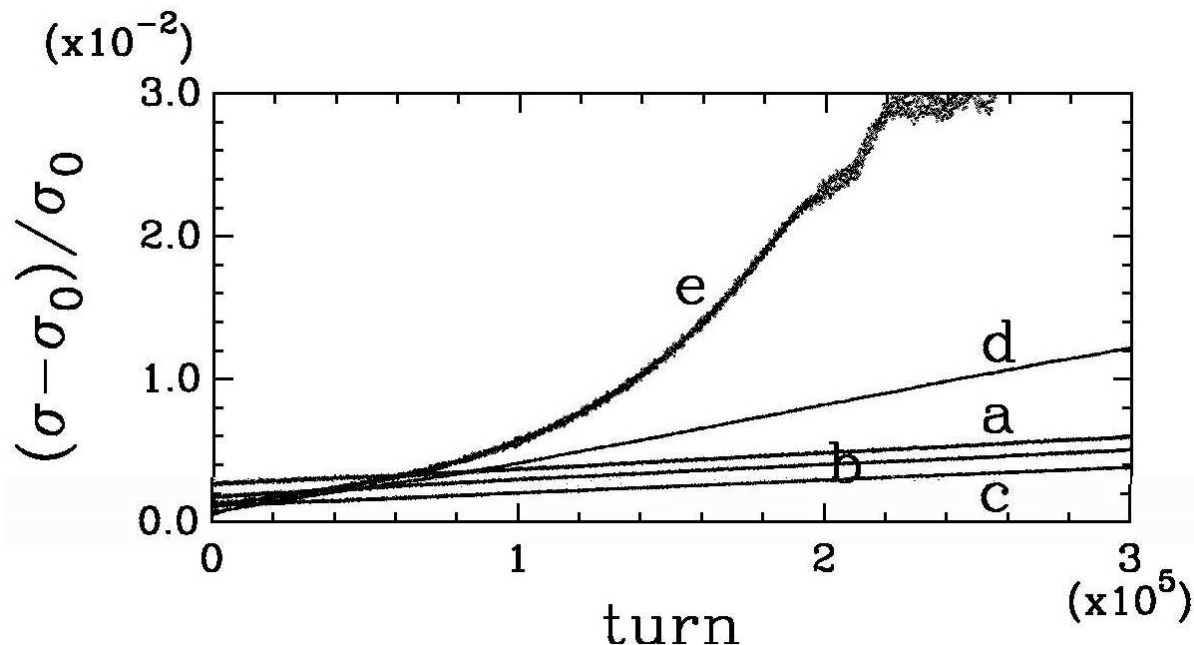




- A reduction of the beam-beam tune spread with electron beams could weaken the Landau damping that is important to the suppression of chaotic coherent beam-beam instability.
- At a “bad” working point, a reduction of the beam-beam tune spread could benefit beams if the weakened Landau damping is still enough to curb chaotic coherent beam-beam instability.

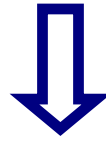
2 π -Cancellation Between \bar{p} -p and \bar{p} -e Collision

Evolution of r.m.s. Beam-size of the Antiproton Beam

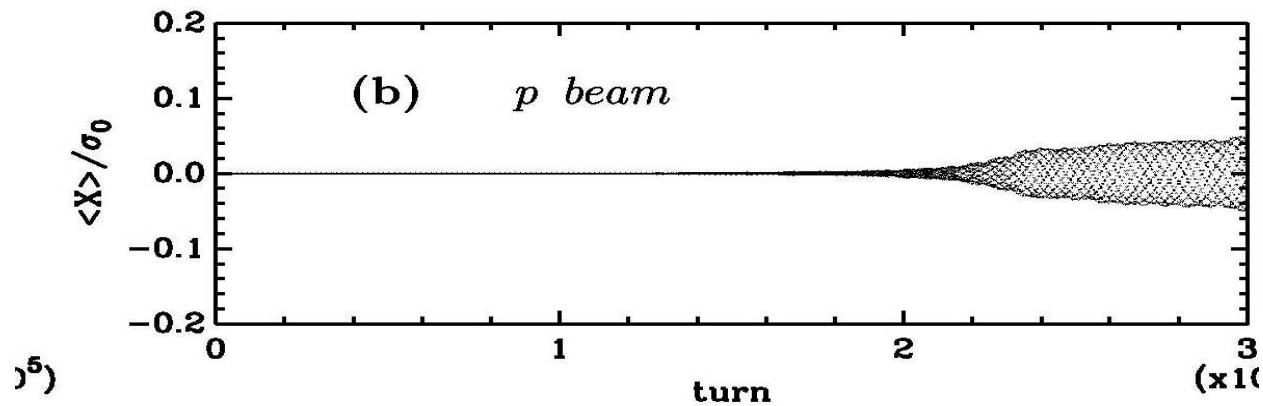
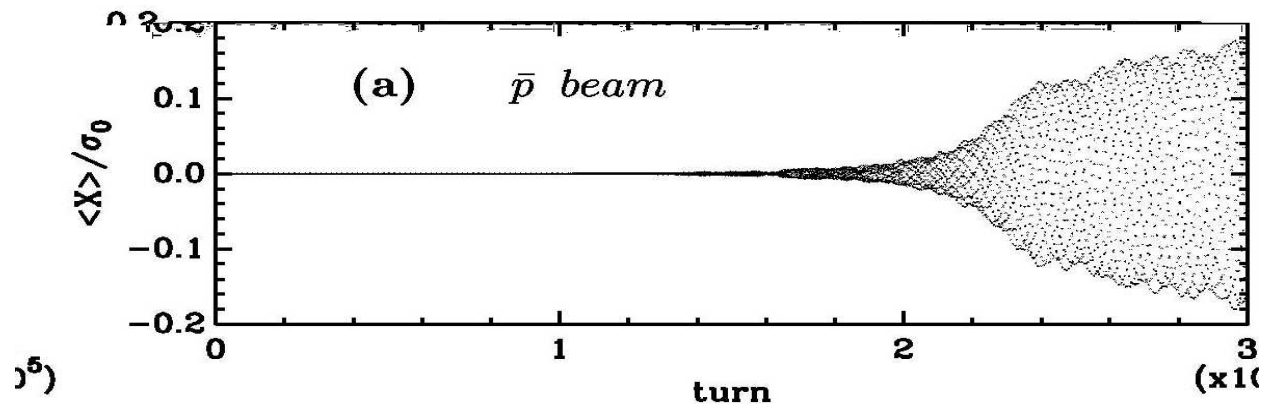


- a: no compensation
- b: 25% compensation
- c: 50% compensation
- d: 75% compensation
- e: Full compensation

Microscopic difference in the distribution of the e and p beam makes otherwise a perfect cancellation of beam-beam interactions failed in the full compensation due to the onset of the chaotic coherent beam-beam instability.



Chaotic Coherent Motion of the Antiproton beam with the full compensation





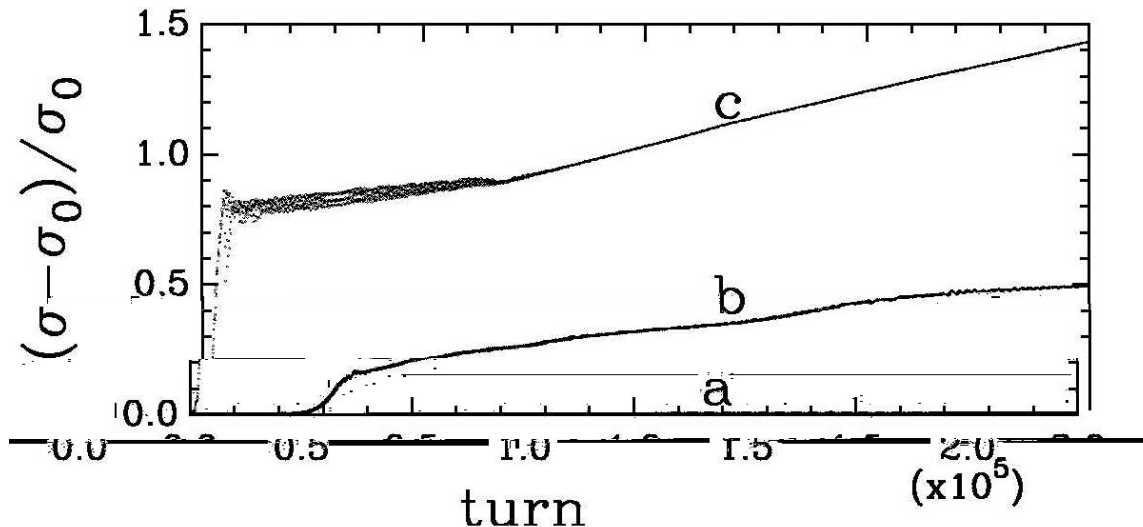
Compensation of Beam-Beam Tune Spread with Electron Beams

LHC MODEL

$$(\nu_x, \nu_y) = (0.31, 0.32)$$

$$\xi_{pp} = 0.01, \text{ 2IPs}$$

Evolution of r.m.s. Beam-size of the proton Beam



a: no compensation

b: 50% compensation

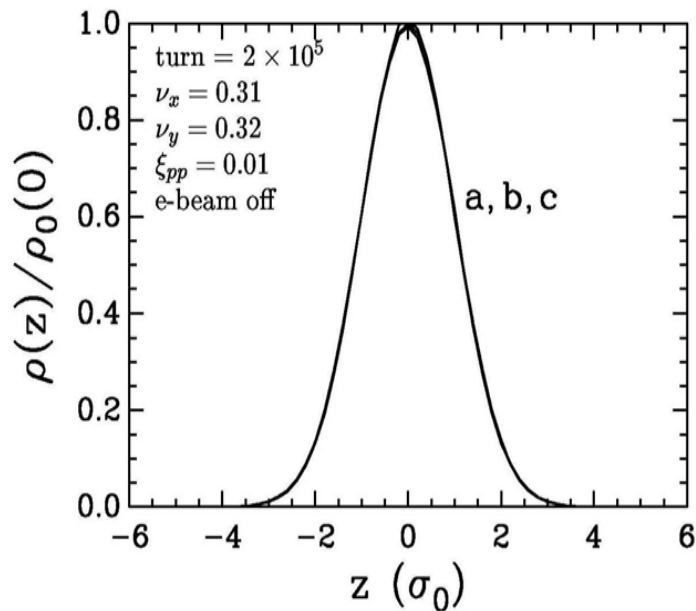
c: Full compensation

Similar to Tevatron, at a “good” working point, the compensation of beam-beam tune spread could damage the beam stability.

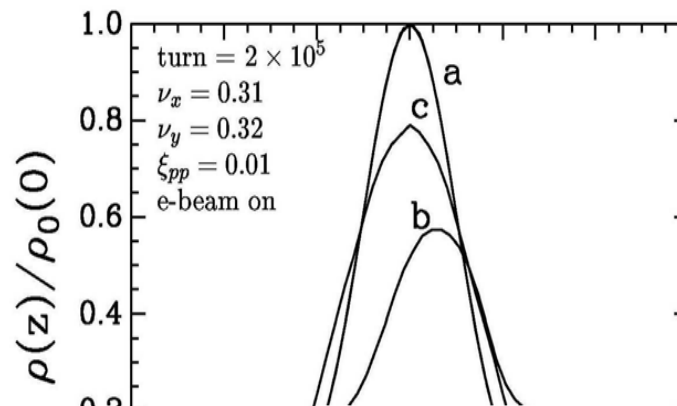


Effects of the Tune-Spread Compensation on Beam Particle Distribution

No Compensation



Full Compensation



a. Initial Gaussian distribution

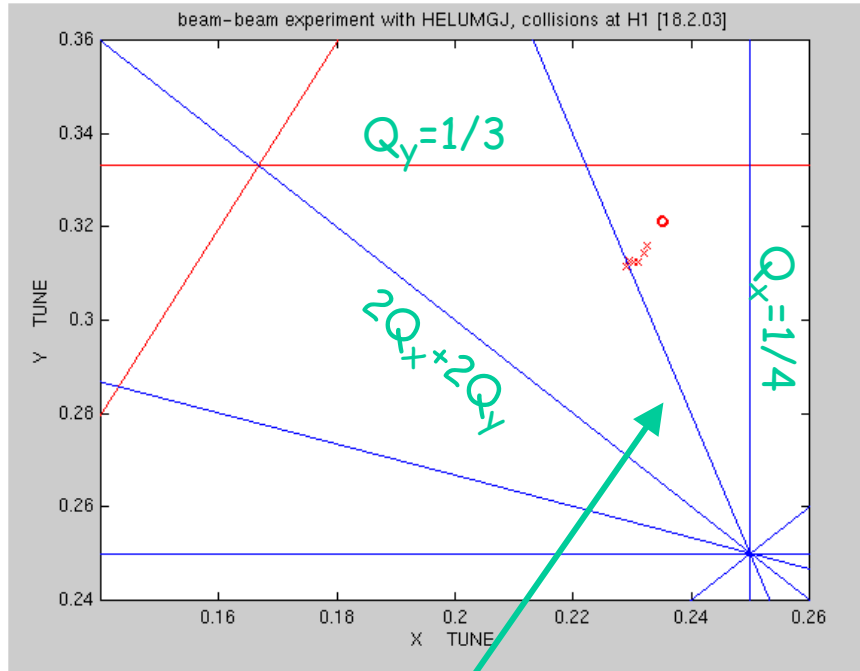
b. Projection in X

c. Projection in Y

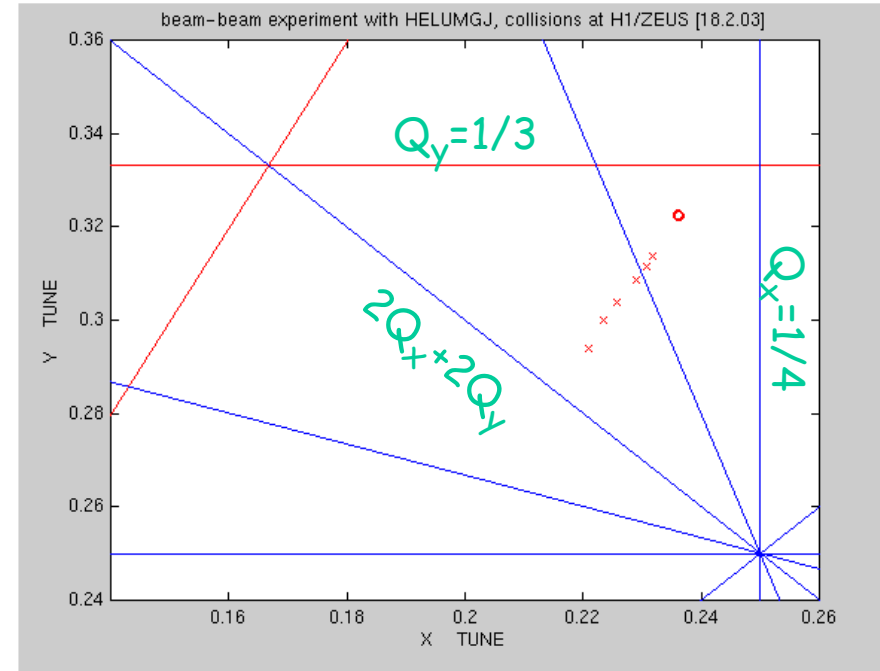
The compensation of beam-beam tune spread could damage beams by inducing the chaotic coherent beam-beam instability that leads to a formation of beam halos.

measured coherent tunes with collisions with 1 and 2 IPs (expmts IA/B)

IA) collisions at H1 only



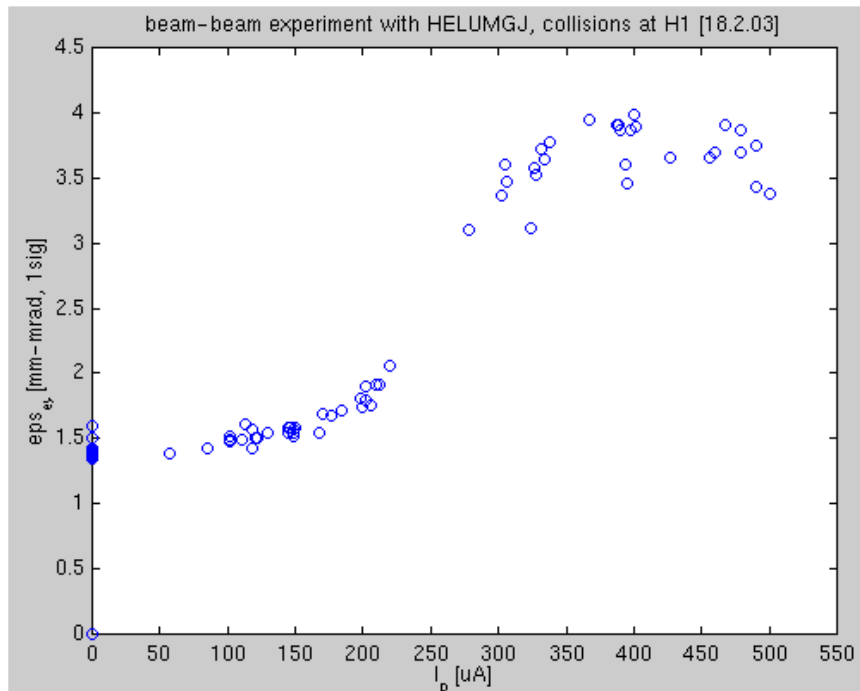
IB) collisions at H1 and ZEUS



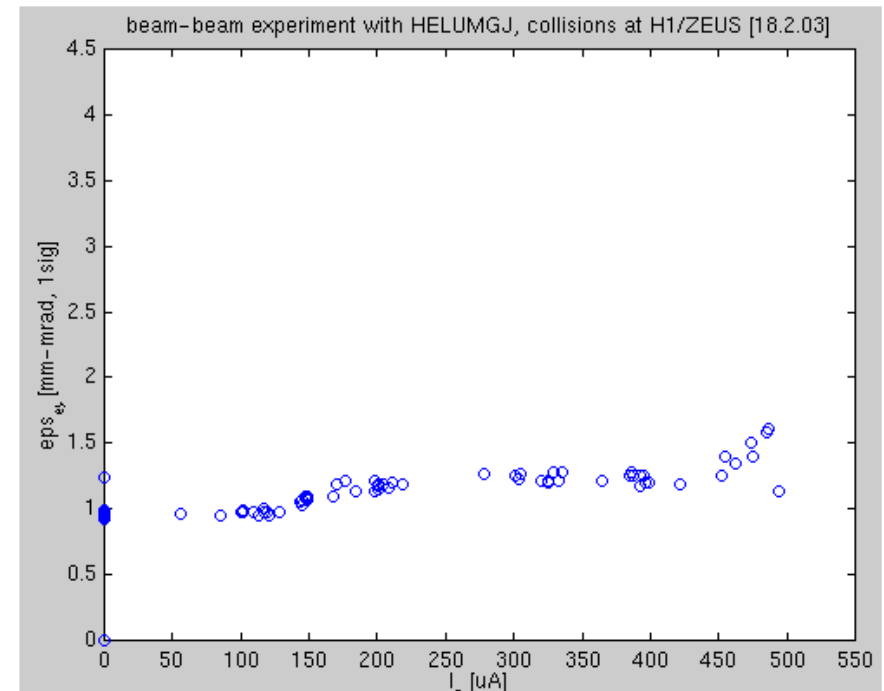
observation: with collisions only at H1, the beam appears to lock onto the $Q_y + 3Q_x$ resonance thereby causing the increased positron vertical emittance; why this is not the case with 2 IPs is unclear

measured positron y emittance vs proton beam current (expmts. IA/B)

IA) collisions at H1 only



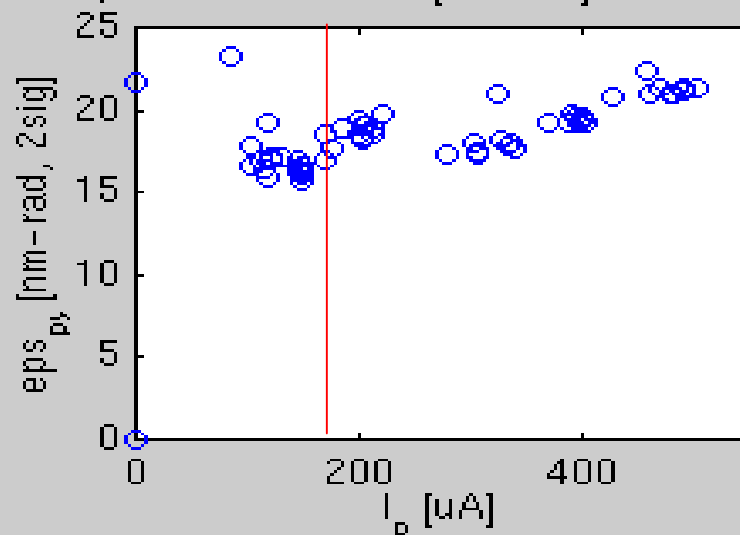
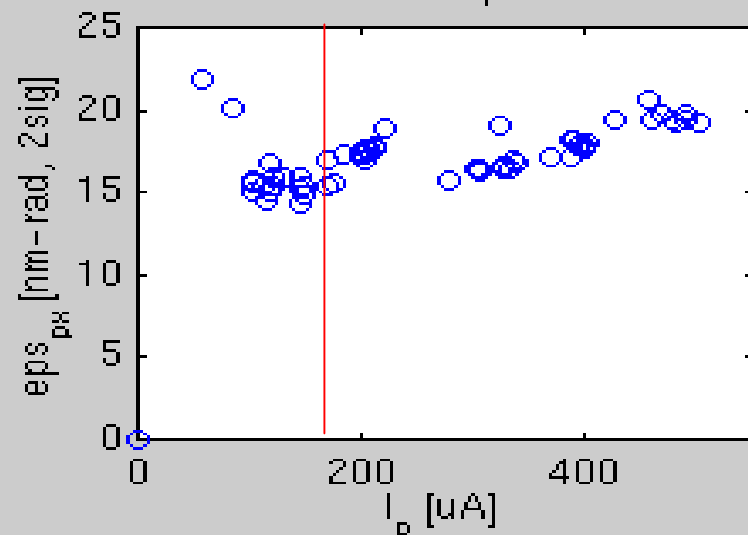
IB) collisions at H1 and ZEUS



observation: significant increase in positron emittance with proton current in experiment with collisions at H1 only

$\epsilon_{p,x}$, $\epsilon_{p,y}$, $\epsilon_{e,y}$ v.s. Proton Bunch Current

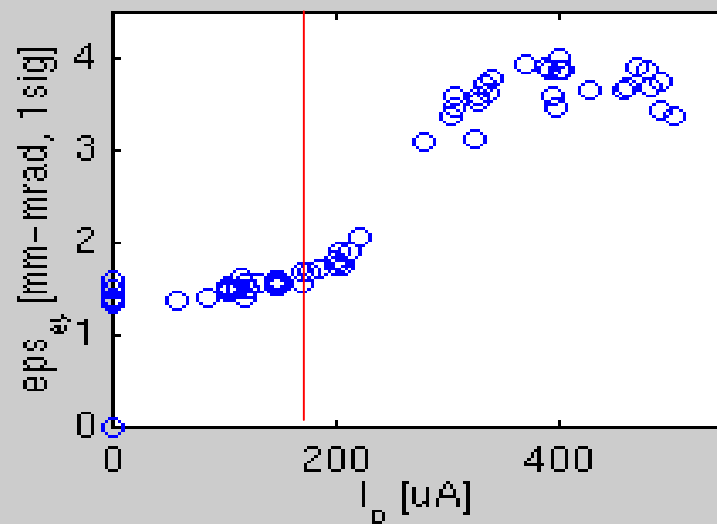
beam-beam experiment with HELUMGJ, collisions at H1 [18.2.03]



$$\xi_c : I_p \approx 170 \mu\text{A}$$

Collective Beam-Beam Instability:

When $\xi_e > \xi_c$ the proton emittance increases with the beam-beam parameter of the positron beam. In this case, $\xi_e \gg \xi_p$.





Summary

評 In a hadron collider with high-intensity beams, the beam-beam interaction is likely to be in the nonlinear regime in which the chaotic coherent beam-beam instability is important. In this situation, having a larger tune spread could be better to the beam stability since it could result in a stronger Landau damping that could suppress the coherent beam-beam instability and, moreover, the existence of a large tune spread reduces the possibility of trapping particles inside a bad resonance.

評 In the case that the working point is close to major resonance, a compensation of the beam-beam tunes spread up to certain degree could improve beam dynamics if the Landau damping is still strong enough for the suppression of the coherent beam-beam instability after the compensation or the damage effects of the nonlinear phase-dependent beam-beam perturbations can be outweighed by the benefit of the tune-spread compensation.